

A semi-implicit NH quasi-elastic soundproof mass-based dynamical system for NWP applications

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- 1 Elaboration of a sound-proof NWP system free from the numerical adverse effects of the fast vertically propagation acoustic waves.
- 2 Maintain the existing AROME/ALARO Dynamical software.

Basic underlying idea :

The actual density $\rho(\theta, p)$ is replaced by a pseudo-density $\tilde{\rho}(\theta) \Rightarrow$ approximate mass-continuity equation:

$$\frac{D\tilde{\rho}}{Dt} = \left(\frac{\partial \tilde{\rho}}{\partial \theta} \right)_p \frac{D\theta}{Dt} + \underbrace{\left(\frac{\partial \tilde{\rho}}{\partial p} \right)_\theta}_{=0} \frac{Dp}{Dt} = -\tilde{\rho} D_3$$

- 1 Quasi-elastic approximation: $\tilde{\rho} = \tilde{\rho}[\theta(x, y, z, t)] \Rightarrow$ the most accurate sound-proof approx. with no evidence of Rossby waves distortion [Arakawa and Konor, 2009]
 \Rightarrow [viable at synoptic and meso-scales](#) .

Basic definitions and approximations

QE reference state

$(\pi, \tilde{\rho}, \tilde{T})$ are thermodynamic variables of the QE reference state defined as

$$\begin{aligned}\tilde{\rho} &= p_{00}^{\kappa} \frac{\pi^{1-\kappa}}{R\theta} \equiv \tilde{\rho}(\pi, \theta) \\ \frac{\partial \pi}{\partial z} &= -\tilde{\rho} g \\ \tilde{T} &= \frac{\pi}{R\tilde{\rho}}\end{aligned}$$

Hence, $\pi = \pi(\theta) \Rightarrow \tilde{\rho} = \tilde{\rho}(\theta)$

Defining pressure departure as $q = \ln(p/\pi) \Rightarrow$

$$p = \pi \exp[q]$$

$$\rho = \tilde{\rho} \exp[(1 - \kappa)q]$$

$$T = \tilde{T} \exp[\kappa q]$$

QE approximation

$$\frac{\dot{\tilde{\rho}}}{\tilde{\rho}} = -D_3 \quad (\text{QE mass-continuity eq.})$$

$$\frac{\dot{\tilde{T}}}{\tilde{T}} = \frac{R}{C_p} \frac{\dot{\tilde{\pi}}}{\tilde{\pi}} + \frac{Q}{C_p T}, \quad (\text{Thermodynamic eq.})$$

QE perfect gaz law $\pi = \tilde{\rho} R \tilde{T} \Rightarrow$

$$\frac{\dot{\tilde{\pi}}}{\tilde{\pi}} + \frac{C_p}{C_v} D_3 = \frac{Q}{C_v T}, \quad (\text{QE momentum constraint eq.})$$

Mass-based vertical coordinate transformation

Choosing π as a vertical coordinate [Laprise 92]

Coordinate transformation $z \rightarrow \pi \rightarrow \eta \in [0, 1] : \pi = A(\eta) + B(\eta)\pi_s$ and $m = \partial\pi/\partial\eta$,

$$\frac{\partial\phi}{\partial\eta} = -\frac{1}{\tilde{\rho}} \frac{\partial\pi}{\partial\eta} = -m \frac{R\tilde{T}}{\pi}$$

[Kasahara] mass continuity equation \Rightarrow

$$\frac{\partial m}{\partial t} + \nabla \cdot (mV) + \frac{\partial}{\partial\eta}(m\dot{\eta}) = 0$$

with $\dot{\eta}(0) = \dot{\eta}(1) = 0 \Rightarrow$

$$\begin{aligned}\dot{\pi} &= V \cdot \nabla\pi - \int_0^\eta \nabla \cdot (mV) d\eta' \\ m\dot{\eta} &= B(\eta) \int_0^1 \nabla \cdot (mV) d\eta' - \int_0^\eta \nabla \cdot (mV) d\eta'\end{aligned}$$

Choosing π as a vertical coordinate [Laprise 92]

Coordinate transformation rules :

$$\frac{\partial}{\partial z} = -g \frac{\tilde{\rho}}{m} \frac{\partial}{\partial \eta}$$

$$\nabla_z = \nabla + \frac{\tilde{\rho}}{m} \nabla \phi \frac{\partial}{\partial \eta}$$

$$\frac{d}{dt} = \frac{\partial}{\partial t} + \mathbf{V} \cdot \nabla + \dot{\eta} \frac{\partial}{\partial \eta}$$

Momentum equations \Rightarrow

$$\frac{dV}{dt} = -\frac{\tilde{\rho}}{\rho} \left[\frac{1}{m} \frac{\partial p}{\partial \eta} \nabla \phi + \frac{1}{\rho} \left(1 + \frac{\rho - \tilde{\rho}}{\tilde{\rho}} \right) \nabla p \right] + \mathcal{F}_V$$

$$\frac{dw}{dt} = g \frac{\tilde{\rho}}{\rho} \left[\frac{1}{m} \frac{\partial (p - \pi)}{\partial \eta} - \frac{\rho - \tilde{\rho}}{\tilde{\rho}} \right] + \mathcal{F}_w$$

Prognostic equations :

$$\frac{dV}{dt} = -e^{\kappa q} \left[\left(1 + \frac{1}{m} \frac{\partial q}{\partial \eta} \right) \nabla \phi + R \tilde{T} \left(\frac{\nabla \pi}{\pi} + \nabla q \right) \right] + \mathcal{F}_V$$

$$\frac{dw}{dt} = \frac{g}{\kappa} \left[\frac{\pi}{m} \frac{\partial (e^{\kappa q} - 1)}{\partial \eta} + \kappa (e^{\kappa q} - 1) \right] + \mathcal{F}_w$$

$$\frac{d\tilde{T}}{dt} = \frac{R\tilde{T}}{C_p} \frac{\dot{\pi}}{\pi} + \frac{Qe^{-\kappa q}}{C_p}$$

$$\frac{\partial \pi_s}{\partial t} = - \int_0^1 \nabla \cdot (mV) d\eta'$$

Diagnostic equations :

$$\tilde{\rho} = \frac{\pi}{R\tilde{T}}$$

$$\nabla\phi = \nabla\phi_s + \int_{\eta}^1 \nabla \left(m \frac{R\tilde{T}}{\pi} \right) d\eta'$$

$$\frac{\dot{\pi}}{\pi} = \mathbf{V} \cdot \frac{\nabla\pi}{\pi} - \frac{1}{\pi} \int_0^{\eta} \nabla \cdot (m\mathbf{V}) d\eta'$$

$$m\dot{\eta} = B(\eta) \int_0^1 \nabla \cdot (m\mathbf{V}) d\eta' - \int_0^{\eta} \nabla \cdot (m\mathbf{V}) d\eta'$$

$$D_3 = \nabla \cdot \mathbf{V} + \frac{\tilde{\rho}}{m} \nabla\phi \cdot \frac{\partial \mathbf{V}}{\partial \eta} - g \frac{\tilde{\rho}}{m} \frac{\partial w}{\partial \eta}$$

QE mass-based dynamical continuous system

QE momentum constraint :

$(V, w, \tilde{T}, \pi_s) \rightarrow$ prognostic variables. $\dot{q} = 0 \Rightarrow q$ is diagnostically determined in such a way the velocity fields satisfy

$$D_3 + \frac{C_v}{C_p} \frac{\dot{\pi}}{\pi} = \frac{Q}{C_p} e^{-\kappa q} \quad \Rightarrow$$

$$D - \frac{C_v}{C_p} \frac{1}{\pi} \int_0^\eta (m D) d\eta' + \overbrace{\left(-g \frac{\tilde{\rho}}{m} \frac{\partial w}{\partial \eta} \right)}^d + \underbrace{\frac{\tilde{\rho}}{m} \nabla \phi \cdot \frac{\partial V}{\partial \eta} + \frac{C_v}{C_p} \left[V \cdot \frac{\nabla \pi}{\pi} - \frac{1}{\pi} \int_0^\eta V \cdot \nabla m d\eta' \right]}_Y = \frac{Q}{C_p} e^{-\kappa q}$$

\Rightarrow

$$\left(1 - \frac{C_v}{C_p} S \right) D + d + Y = \frac{Q}{C_p} e^{-\kappa q}$$

QE mass-based dynamical continuous system

Use of $dl = d + Y$ as prognostic variable in lieu of w

Vertical momentum eq. \Rightarrow

$$\frac{ddl}{dt} = -\frac{g^2}{R\tilde{T}} \frac{\pi}{m} \frac{\partial}{\partial \eta} \left[\left(\frac{\pi}{m} \frac{\partial}{\partial \eta} + \kappa \right) \left(\frac{e^{\kappa q} - 1}{\kappa} \right) + \frac{\mathcal{F}_w}{g} \right] + \left(\frac{\pi}{m} \frac{\partial V}{\partial \eta} \right) \cdot \frac{g \nabla w}{R\tilde{T}} + (dl - Y)(D - D_3) + \dot{Y},$$

QE momentum constraint \Rightarrow

$$\left(1 - \frac{C_v}{C_p} S \right) D + dl = \frac{Q}{C_p} e^{-\kappa q}$$

Diagnostic for $w \Rightarrow$

$$w = w_s + \int_{\eta}^1 m \frac{R\tilde{T}}{g\pi} (dl - Y) d\eta'$$

Basic concept

Defining $X = [V, w, \tilde{T}, \tilde{q}_S]$ the vector of prognostic variables and $Z = [X, q]$ the state vector, the QE system can be symbolically written as

$$\frac{\partial X}{\partial t} = \mathcal{A}(X) + \hat{\mathcal{M}}_x(Z) \equiv \mathcal{M}(Z)$$

together with the QE momentum constraint :

$$\hat{\mathcal{M}}_q(X) = 0$$

The physical contributions are not taken into account.

A proposed SI time discretization

Basic concept

As a classical 3-TL SI Eulerienne approach :

$$\frac{X^+ - X^-}{2\Delta t} = \mathcal{M}(Z^0) - \mathcal{L}^* \cdot Z^0 + \frac{\mathcal{L}^* \cdot Z^- + \mathcal{L}^* \cdot Z^+}{2},$$

with

$$\hat{\mathcal{M}}_q(X^+) = 0$$

The prognostic fields discretely satisfy the non-linear QE momentum constraint at $t + \Delta t$.

A proposed SI time discretization

Basic concept

The SI scheme is combined with a Newton-Raphson-like iterative treatment of the QE momentum constraint, it yields

$$\frac{X^{+(\nu)} - X^-}{2\Delta t} = \mathcal{M}(Z^0) - \mathcal{L}^*.Z^0 + \frac{\mathcal{L}^*.Z^- + \mathcal{L}^*.Z^{+(\nu)}}{2},$$

with

$$\hat{\mathcal{M}}_q [X^{+(\nu-1)}] - \mathcal{C}^*.X^{+(\nu-1)} + \mathcal{C}^*.X^{+(\nu)} = 0$$

- 1 Stability is controlled by the magnitude of the explicit residual terms $(\mathcal{M} - \mathcal{L}^*)$.
- 2 Convergence is determined by the magnitude of the non-linear residual term $(\hat{\mathcal{M}}_q - \mathcal{C}^*)$ of QE momentum constraint.

Iterative linear implicit problem

For $\nu \in [1, N_{\text{stop}}]$,

$$X^{+(\nu)} - \Delta t \mathcal{L}^* . Z^{+(\nu)} = X^\bullet$$

$$\mathcal{C}^* . X^{+(\nu)} = \mathcal{R}^{+(\nu-1)}$$

- 1 Our choice of \mathcal{L}^* and \mathcal{C}^* consists in linearizing the full dynamical system around a resting stationary, isothermal and horizontally homogeneous reference-state X^\bullet
- 2 N_{stop} is determined in such a way that the QE constraint $\hat{\mathcal{M}}_q(Z^{+(N_{\text{stop}})}) \simeq 0$ is fulfilled at an acceptable degree of accuracy (machine precision is typically required).

Iterative linear implicit problem

$$D^{+(\nu)} + \Delta t RT^* \nabla^2 \left[\mathbf{G}^* \tilde{T}_*^{+(\nu)} + \tilde{q}_S^{+(\nu)} + q^{+(\nu)} \right] = \nabla \cdot \mathcal{V}^\bullet$$

$$dl^{+(\nu)} + \Delta t \frac{g}{rH_*} \mathbf{L}_\kappa^* q^{+(\nu)} = \mathcal{D}^\bullet,$$

$$\tilde{T}_*^{+(\nu)} + \Delta t \kappa \mathbf{S}^* D^{+(\nu)} = \mathcal{T}^\bullet,$$

$$\tilde{q}_S^{+(\nu)} + \Delta t \mathbf{N}^* D^{+(\nu)} = \mathcal{P}_S^\bullet$$

with

$$\mathbf{S}_\kappa^* D^{+(\nu)} + dl^{+(\nu)} = \mathcal{R}^{+(\nu-1)} \equiv (\mathbf{S}_\kappa^* - \mathbf{S}_\kappa) D^{+(\nu-1)},$$

Algebraical constraints on the vertically discrete linear model

Considering \mathbf{G}^* , \mathbf{S}^* , \mathbf{N}^* , and \mathbf{L}_v^* as already existing operators for EE system,

- 1 $\mathbf{S}_{\kappa}^* = [\mathbf{I} - (1 - \kappa)\mathbf{S}]$ is an invertible integral matrix.
- 2 $\mathbf{L}_{\kappa}^* = \mathbf{S}_{\kappa}^* \mathbf{L}_v^*$ is akin as a tridiagonal vertical Laplacian operator
- 3 The following identity should be satisfied :

$$\mathbf{L}_v^*[\mathbf{G}^* \mathbf{S}^* + (1/\kappa)\mathbf{N}^*] = -\mathbf{T}^*$$

where \mathbf{T}^* is a tridiagonal vertical averaging operator similar to the one describes in BHBG95. hereafter $\mathbf{G}^* \mathbf{S}^* + (1/\kappa)\mathbf{N}^* \equiv \mathbf{K}^*$.

Resolution of the iterative linear implicit problem

Elimination procedure

$$[\mathbf{I} - \Delta t^2 N_*^2 \mathbf{K}^* \nabla_*^2] D^{+(\nu)} + \frac{1}{\tau^*} \nabla_*^2 q^{+(\nu)} = D^\bullet,$$

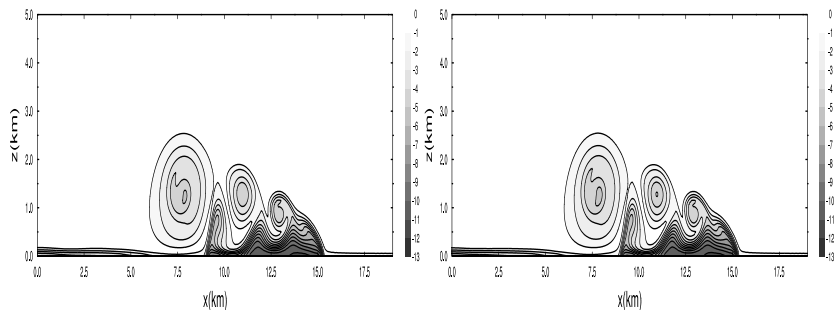
$$D^{+(\nu)} - \frac{1}{\tau^* r} \mathbf{L}_\nu^* q^{+(\nu)} = \mathbf{S}_\kappa^{*-1} [\mathcal{R}^{+(\nu-1)} - \mathcal{D}^\bullet],$$

\Rightarrow For $\nu \in [1, N_{\text{stop}}]$

$$[\mathbf{L}_\nu^* + (r\mathbf{I} + \Delta t^2 N_*^2 \mathbf{T}^*) \nabla_*^2] D^{+(\nu)} = D^{\bullet\bullet} \\ + r \nabla_*^2 \mathbf{S}_\kappa^{*-1} (\mathbf{S}_\kappa^* - \mathbf{S}_\kappa) D^{+(\nu-1)},$$

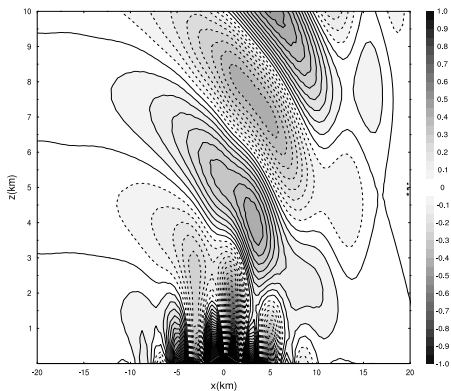
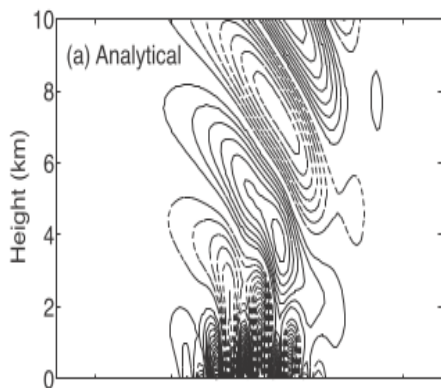
Idealized 2D test-case

Density current test case: Distributions of perturbations of potential temperature for the cold bubble experiment : spacing $\Delta x = 50$ m and $\Delta z \approx 50$ m, with $\Delta t = 0.25$ s. EE solution (left) and QE solution (right) at $t = 900$ s



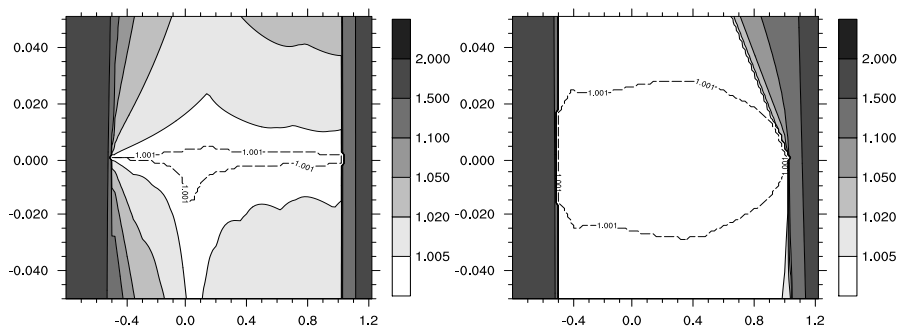
Idealized 2D test-case

Schär mountain test case : Numerical solution of vertical velocity (m s^{-1}) at $t = 12$ h, with $\Delta t = 4$ s, $\Delta x = 500$ m, and $\Delta z \simeq 300$ m.



SHB linear stability analysis

stability as function a uniform slope s (%) and the thermal non-linear residual factor $\alpha = (T - T^*)/T^*$ actual 3-TL SI EE (left) 3-TL SI QE (right)



Summary and perspectives

- 1 QE soundproof system can be readily blended into the existing HPE/EE ARPEGE/AROME Dynamical kernel using Laprise approach.
- 2 A semi-implicit time discretization of such a soundproof system is now possible in a similar framework as in AROME (extension to ICI scheme is also possible).
- 3 Numerical testing with SI QE system seems to present a better stability property than SI EE system while maintaining a similar degree of accuracy.